

CH 7: Connections

7.1 Affine Connections

Dfn 7.1: Affine connection on M is an \mathbb{R} -Bilinear map

$$\nabla: \mathfrak{X}(M) \times \mathfrak{X}(M) \rightarrow \mathfrak{X}(M), \quad (X, Y) \mapsto \nabla_X Y$$

satisfying $\forall X, Y \in \mathfrak{X}(M)$ and $f \in C^\infty(M)$

(i) $\nabla_{fX} Y = f \nabla_X Y$

(ii) $\nabla_X (fY) = X(f)Y + f \nabla_X Y$

\mathbb{R} -bilinearity says: $\nabla_{X_1+X_2} Y = \nabla_{X_1} Y + \nabla_{X_2} Y$
 $\nabla_X (Y_1+Y_2) = \nabla_X Y_1 + \nabla_X Y_2$

Abuse of notation: $\partial_i = \frac{\partial}{\partial x^i}$ and $\nabla_i = \nabla_{\partial_i}$

Connection coefficients: $\nabla_i \partial_j = \Gamma_{ij}^k \partial_k$ using Einstein summation convention.

Local expression: $\nabla_X Y = X^i \partial_i Y^j \partial_j + X^i Y^j \Gamma_{ij}^k \partial_k$

Difference of two affine connections is tensorial:

Lem 7.2: ∇ and ∇' affine connections. The following is a $(1,2)$ -tensor field:

$$\kappa: \mathfrak{X}(M) \times \mathfrak{X}(M) \rightarrow \mathfrak{X}(M) \quad \kappa: (X, Y) \mapsto \nabla_X Y - \nabla'_X Y$$

b.c. κ defines a trilinear map

$$\kappa: \Omega^1(M) \times \mathfrak{X}(M) \times \mathfrak{X}(M) \rightarrow C^\infty(M), \quad \kappa(\alpha, X, Y) = \alpha(\kappa(X, Y))$$

Standard affine connection on \mathbb{R}^n : $\Gamma_{ij}^k \equiv 0$.

Dfn 7.3: Parallel vector fields: a v.f. $Y \in \mathfrak{X}(M)$ such that $\nabla_X Y = 0 \quad \forall X \in \mathfrak{X}(M)$ is called parallel.

Lem 7.5: Let ∇' and ∇'' be affine connections on M . Let $p_1, p_2 \in C^\infty(M)$ with $p_1 + p_2 = 1$. Then $\nabla = p_1 \nabla' + p_2 \nabla''$ is an affine connection.

Prop 7.6: $\forall X, Y \in \mathfrak{X}(M)$ define $\nabla: \mathfrak{X}(M) \times \mathfrak{X}(M) \rightarrow \mathfrak{X}(M)$ by $\nabla_X Y := \sum_{i \in \mathbb{Z}} p_i \nabla_X^{(i)} Y$.

7.2 : Torsion and Curvature

Dfn 7.7: Torsion of ∇ : $T^\nabla: \mathfrak{X}(M) \times \mathfrak{X}(M) \rightarrow \mathfrak{X}(M)$ (1,2)-tensor field
 $T^\nabla(X, Y) := \nabla_X Y - \nabla_Y X - [X, Y]$ prop 7.8

Curvature of ∇ : $R^\nabla: \mathfrak{X}(M) \times \mathfrak{X}(M) \times \mathfrak{X}(M) \rightarrow \mathfrak{X}(M)$ (1,3)-tensor field
 $R^\nabla(X, Y)Z := \nabla_X \nabla_Y Z - \nabla_Y \nabla_X Z - \nabla_{[X, Y]} Z$

"flat" = zero curvature.
 "torsion-free" = zero torsion.

Torsion coefficients: $T_{ij}^k \partial_k = T(\partial_i, \partial_j)$
 $\Rightarrow T_{ij}^k = \Gamma_{ij}^k - \Gamma_{ji}^k$

Curvature coefficients: $R_{ijk}^l \partial_l = R^\nabla(\partial_i, \partial_j)\partial_k$
 $\Rightarrow R_{ijk}^l = \partial_i \Gamma_{jk}^l - \partial_j \Gamma_{ik}^l + \Gamma_{jk}^m \Gamma_{im}^l - \Gamma_{ik}^m \Gamma_{jm}^l$

Dfn 7.2 : Ricci tensor: (0,2)-tensor defined by

$Ric^\nabla(X, Y) := \text{tr}(Z \rightarrow R(Z, X)Y)$
 $R_{ij} := Ric(\partial_i, \partial_j) \Rightarrow R_{ij} = R_{kij}^k$ } Ric is a contraction of R.

7.3 : Bianchi Identities

Extending by Leibniz rule:

- $f \in C^\infty(M)$, define $\nabla_X(f) := X(f)$
- $\alpha \in \Omega^1(M)$, define $\nabla_X(\alpha(Y)) = (\nabla_X \alpha)(Y) + \alpha(\nabla_X Y)$

Define $\mathcal{G}A(X, Y, Z) := A(X, Y, Z) + A(Y, Z, X) + A(Z, X, Y)$

Prop 7.6: Bianchi identity:

1. Algebraic Bianchi identity:

$\mathcal{G}R(X, Y)Z = \mathcal{G}(\nabla_X T)(Y, Z) + \mathcal{G}T(T(X, Y), Z)$

2. Differential Bianchi identity:

$\mathcal{G}((\nabla_Z R)(X, Y))W + \mathcal{G}R(T(X, Y), Z)W = 0$

7.4 Koszul Connections

Dfn 7.18: A **Koszul connection** on a vector bundle $E \rightarrow M$ is an \mathbb{R} -bilinear map

$$\nabla : \mathcal{X}(M) \times \Gamma(E) \rightarrow \Gamma(E), \quad (X, s) \mapsto \nabla_X s$$

Satisfying $\forall f \in C^\infty(M), X \in \mathcal{X}(M)$ and $s \in \Gamma(E)$:

- (i) $\nabla_{fX} s = f \nabla_X s$, and
- (ii) $\nabla_X (fs) = X(f)s + f \nabla_X s$

e.g. affine connection = Koszul connection on tangent bundle $TM \rightarrow M$.

Let ∇^E and ∇^F be Koszul connections on $E \rightarrow M$ and $F \rightarrow M$ respectively. Induces a Koszul connection ∇^\otimes on the vector bundle $E \otimes F \rightarrow M$ as follows:

$$\nabla_X^{E \otimes F} (s \otimes t) := (\nabla_X^E s) \otimes t + s \otimes (\nabla_X^F t)$$

and ∇^{Hom} on the vector bundle $\text{Hom}(E, F) \rightarrow M$ as follows:

$$(\nabla_X^{\text{Hom}} \varphi)(s) := \nabla_X^F (\varphi(s)) - \varphi(\nabla_X^E s).$$

Curvature of Koszul connection: $R^\nabla(X, Y)s = \nabla_X \nabla_Y s - \nabla_Y \nabla_X s - \nabla_{[X, Y]} s.$

Connection coefficients (for rank- m v.b) are mn^2 functions $\omega_i^b{}_a \in C^\infty(U)$ defined by

$$\nabla_i e_a := \omega_i^b{}_a e_b$$

where e_a = local sections linearly independent on U .

$$\begin{aligned} \text{Local expression of } s = s^a e_a &\Rightarrow \nabla_i s = \partial_i s^a e_a + s^a \omega_i^b{}_a e_b \\ &= (\partial_i s^a + \omega_i^a{}_b s^b) e_a \end{aligned}$$

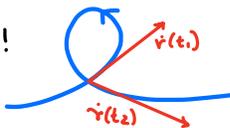
7.5: Parallel Transport:

Smooth curve: a smooth map $\gamma: [0,1] \rightarrow M$

Section along γ : smooth map $\sigma: [0,1] \rightarrow E$ such that $\pi \circ \sigma = \gamma$, equivalently $\sigma(t) \in E_{\gamma(t)}$.

\hookrightarrow if $\sigma \in \Gamma(E)$, then $\sigma = s \circ \gamma$ is a section along γ for **extendible sections**.

NB: not all sections are extendible!!



$\Gamma_{\gamma}(E) =$ space of sections along $\gamma = C^{\infty}([0,1])$ -module: $(f\sigma)(t) = f(t)\sigma(t) \in E_{\gamma(t)}$.

any section $\sigma \in \Gamma_{\gamma}(E)$ can be written locally as $\sigma = \sigma^a(e_a \circ \gamma)$.

Dfn 7.24: ∇ Koszul on $E \rightarrow M$. The **covariant derivative** on $\Gamma_{\gamma}(E)$ induced by ∇ is the **\mathbb{R} -linear map**

$$\frac{D}{dt} : \Gamma_{\gamma}(E) \rightarrow \Gamma_{\gamma}(E)$$

characterized by the properties that:

$$(i) \forall s \in \Gamma(E), \frac{D}{dt}(s \circ \gamma) = (\nabla_{\dot{\gamma}} s) \circ \gamma$$

$$(ii) \forall f \in C^{\infty}([0,1]) \text{ and } \sigma \in \Gamma_{\gamma}(E) \quad \frac{D}{dt}(f\sigma) = \frac{df}{dt}\sigma + f \frac{D\sigma}{dt}$$

where $\dot{\gamma} = \gamma_{*} \left(\frac{d}{dt} \right)$

parallel/covariantly constant: $\frac{D\sigma}{dt} = 0$. } parallel iff $\sigma^a \in C^{\infty}([0,1])$ obeys ODE $\frac{d\sigma^a}{dt} + \dot{\gamma}^i(t) \omega_i^a{}_b(\gamma(t)) \sigma^b(t) = 0$.

Parallel transport: along γ : given $v \in E_{\gamma(0)}$, $\exists!$ $\sigma \in \Gamma_{\gamma}(E)$, $\sigma(0) = v$. Evaluation at 1 gives linear map

$$P_{\gamma} : E_{\gamma(0)} \rightarrow E_{\gamma(1)} \quad v \mapsto \sigma(1).$$

Let $s \in \Gamma(E)$ be s.t. $\nabla_X s = 0 \quad \forall X \in \mathfrak{X}(M)$. Given any curve $\gamma: [0,1] \rightarrow M$, $s \circ \gamma$ is a parallel section along γ .

Prop 7.26: $s_1, s_2 \in \Gamma(E)$ s.t. $\nabla_X s_1 = \nabla_X s_2 = 0 \quad \forall X \in \mathfrak{X}(M)$ and $s_1(p) = s_2(p)$ for some $p \in M$. Then $s_1 = s_2$.

7.6 : Geodesics

Dfn 7.29: ∇ an affine connection on M . A smooth curve $\gamma: I \rightarrow M$ with $I \subset \mathbb{R}$ a nontrivial interval is called a **geodesic** for ∇ if its velocity $\dot{\gamma} := \gamma_* \left(\frac{d}{dt} \right)$, a vector field along γ , is "self-parallel":

$$\frac{D}{dt} \dot{\gamma} = 0.$$

γ is a geodesic iff $\gamma^i(t)$ (where $\dot{\gamma} = \dot{\gamma}^i \partial_i$) satisfy the ODE:

$$\ddot{\gamma}^k(t) + \Gamma_{ij}^k(\gamma(t)) \dot{\gamma}^i(t) \dot{\gamma}^j(t) = 0.$$

notice that torsion does not contribute to eqn since second term is symmetric in $i \leftrightarrow j$.

Prop 7.31: $p \in M$ and $v \in T_p M$. $\exists \epsilon > 0$ and $\exists!$ geodesic $\gamma: [0, \epsilon) \rightarrow M$ with $\gamma(0) = p$ and $\dot{\gamma}(0) = v$.

notation: $\gamma_{p,v} :=$ unique geodesic with ics $\gamma(0) = p$ and $\dot{\gamma}(0) = v$.

Cor 7.32: $\forall p \in M, v \in T_p M$, and $s > 0$, $\exists \epsilon > 0$ such that $\gamma_{p,sv}(t) = \gamma_{p,v}(st)$ for $t \in [0, \epsilon)$.

Dfn 7.33: (**Exponential map**) ∇ affine connection on M and $p \in M$. Exponential map $\exp_p: T_p M \rightarrow M$ is defined to be $\exp_p(v) = \gamma_{p,v}(1)$.

Prop 7.34: In **normal coordinates** at $p \in M$, the connection coefficients satisfy

$$\Gamma_{ij}^k(p) + \Gamma_{ji}^k(p) = 0$$

Prop 7.35: $\gamma: I \rightarrow M$ a geodesic of affine connection ∇ . Let $\phi: M \rightarrow M$ be a diffeomorphism. Then $\gamma^\phi: I \rightarrow M$, defined by $\gamma^\phi = \phi \circ \gamma$, is a geodesic of the transformed connection ∇^ϕ .

\Rightarrow Geodesic can also be described by the general equation

$$\frac{D}{dt} \dot{\gamma} = g(t) \dot{\gamma}$$

for some function g . Affine parameter t for a geodesic is one relative to which the velocity is self-parallel, so that $g(t) = 0$. Two affine parameters s and t are then related by $t = as + b$ for $a > 0, b \in \mathbb{R}$.

CH 8 Riemannian Geometry

8.1 The Metric Tensor

8.1.1 Inner Products

Inner product: a non-degenerate symmetric bilinear form $\langle -, - \rangle : V \times V \rightarrow \mathbb{R}$

$$\langle v, w \rangle = \langle w, v \rangle \quad \langle v_1 + \alpha v_2, w \rangle = \langle v_1, w \rangle + \alpha \langle v_2, w \rangle$$

Non-degenerate: $\langle v, w \rangle = 0 \quad \forall w \in V$, then $v = 0$.

Inner product space = real vector space with an inner product

Isometry: = isomorphism $\phi: V \rightarrow W$ between inner product spaces such that $\forall v_1, v_2 \in V$,

$$\langle \phi(v_1), \phi(v_2) \rangle_W = \langle v_1, v_2 \rangle_V$$

$\mathbb{E}^{s,t} := (\mathbb{R}^n, \langle -, - \rangle)$ is the vector space \mathbb{R}^n with the inner product

$$\langle (v_1, \dots, v_n), (w_1, \dots, w_n) \rangle = \sum_{i=1}^s v_i w_i - \sum_{j=s+1}^n v_j w_j$$

Orthonormal: basis (e_1, \dots, e_n) satisfying $\langle e_i, e_j \rangle = \pm \delta_{ij}$.

Lemma 8.2: Every n -dimensional inner product space $(V, \langle -, - \rangle)$ is isometric to $\mathbb{E}^{s,t}$ for some s and t with $s+t=n$.

Sylvester's Law of Inertia: signature (s,t) is uniquely determined by inner product.

Euclidean vector space: $(n,0)$, $\eta = \begin{pmatrix} 1 & & & \\ & \ddots & & \\ & & 1 & \\ & & & \ddots & \\ & & & & 0 & \\ & & & & & \ddots & \\ & & & & & & 0 & \\ & & & & & & & \ddots & \\ & & & & & & & & 0 & \\ & & & & & & & & & \ddots & \\ & & & & & & & & & & 0 \end{pmatrix}$

Lorentzian vector space: $(n-1,1)$, $\eta = \begin{pmatrix} 1 & & & \\ & \ddots & & \\ & & 1 & \\ & & & \ddots & \\ & & & & 0 & \\ & & & & & \ddots & \\ & & & & & & 0 & \\ & & & & & & & \ddots & \\ & & & & & & & & 0 & \\ & & & & & & & & & \ddots & \\ & & & & & & & & & & 0 \end{pmatrix}$

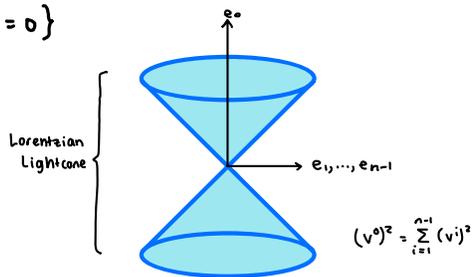
Light cone: $\mathbb{L} = \{ v \in V \mid \langle v, v \rangle = 0 \}$

Nonzero vector is said to be:

Spacelike: $\langle v, v \rangle > 0$

Lightlike/null: $\langle v, v \rangle = 0$

Timelike: $\langle v, v \rangle < 0$



$\begin{cases} \text{future} = \text{timelike and } v^0 > 0 \\ \text{past} = \text{timelike and } v^0 < 0 \\ \vdots \\ \text{spacelike separated from the origin (spacelike vectors).} \end{cases}$

Notation: Symmetric product $uv := \frac{1}{2} (u \otimes v + v \otimes u)$.

8.1.2 Metrics

Dfn 8.3 (Metric): A metric of signature (s,t) with $s+t=n$ on an n -dimensional manifold M is a section $g \in \Gamma(\mathcal{O}^2 T^*M)$ such that for all $p \in M$, g_p defines an inner product on $T_p M$ of signature (s,t) .

Local expression: $g = g_{ij} dx^i dx^j$ where $g_{ij} = g(\partial_i, \partial_j) \in C^\infty(U)$.
 $\forall p \in U$, $[g_{ij}(p)]$ is invertible.

$$g(a,b) = a^T [g] b$$

Riemannian Manifold: (M,g) = manifold equipped with a metric.

Lorentzian Manifold: (M,g) with signature $(n-1,1)$.

Dfn 8.7: $(M,g), (N,h)$ be two Riemannian manifolds. A diffeomorphism $F: M \rightarrow N$ is called an **isometry** if $F^*h = g$; that is, if $\forall p \in M$ and $X_p, Y_p \in T_p M$,

$$g_p(X_p, Y_p) = h_{F(p)}((F_*)_p X_p, (F_*)_p Y_p).$$

Local isometry at $a \in M$ if there's a neighbourhood $U \subset M$ of a s.t. $F|_U: U \rightarrow F(U)$ is an isometry.
 \hookrightarrow notion of (M,g) and (N,h) being locally isometric.

Isometries form a subgroup of the group of diffeomorphisms

Dfn 8.8: A vector field $X \in \mathfrak{X}(M)$ on a Riemannian manifold is called a **"Killing vector field"** if $\mathcal{L}_X g = 0$. That is, if for all $Y, Z \in \mathfrak{X}(M)$,

$$Xg(Y,Z) = g([X,Y], Z) + g(Y, [X,Z]).$$

note: Lie Bracket of two Killing vectors is a Killing vector

(M,g) a Riemannian manifold and $F: N \rightarrow M$ an embedding. The pullback F^*g defines a metric on N .
 \hookrightarrow An embedding $F: N \rightarrow M$ is called isometric if $h = F^*g$.

Example 8.9 = round metric on S^2

Example 8.11 = Minkowski Spacetime, $M = \mathbb{R}^4$, (t, x, y, z) and metric $\eta = -dt^2 + dx^2 + dy^2 + dz^2$.

8.2 The Levi-Civita Connection

Thm 8.12 (Fundamental Theorem of Riemannian Geometry): For a Riemannian manifold (M, g) ,
 $\exists!$ torsion-free affine connection ∇ which is compatible with the metric: $\nabla g = 0$.

Koszul formula for ∇ : $2g(\nabla_X Y, Z) = Xg(Y, Z) + Yg(X, Z) - Zg(X, Y) + g(Z, [X, Y]) - g(Y, [X, Z]) - g(X, [Y, Z])$.

\hookrightarrow ∇ is known as the "Levi-Civita" connection.

Christoffel symbols: connection coefficients relative to a local chart of the Levi-Civita connection.

$$\Gamma_{ij}^m = \frac{1}{2} g^{km} (\partial_i g_{jk} + \partial_j g_{ik} - \partial_k g_{ij})$$

(notice $\Gamma_{ij}^m = \Gamma_{ji}^m$).

\Rightarrow Christoffel symbols vanish in a local chart if the metric coefficients are constant in that chart.

Prop 8.16: Let X be a Killing vector ($\mathcal{L}_X g = 0$). Then X obeys the Killing equation:
 $\forall Y, Z \in \mathfrak{X}(M)$,

$$g(\nabla_Y X, Z) + g(Y, \nabla_Z X) = 0$$

Prop 8.17: $\forall u, v \in \Gamma_\gamma(TM)$ along γ

$$\frac{d}{dt} g(u, v) = g\left(\frac{du}{dt}, v\right) + g\left(u, \frac{dv}{dt}\right).$$

Cor 8.18: $\gamma: I \rightarrow M$ a geodesic for the Levi-Civita connection of a Riemannian manifold (M, g) .
 Then $g(\dot{\gamma}, \dot{\gamma})$ is constant along γ .

Timelike: $g(\dot{\gamma}, \dot{\gamma}) < 0$
 Lightlike or null: $g(\dot{\gamma}, \dot{\gamma}) = 0$
 Spacelike: $g(\dot{\gamma}, \dot{\gamma}) > 0$

} $\leq 0 = \text{causal}$

Ex. 8.19: Minkowski spacetime: geodesic eqn has linear components relative to an affine parameter.

$\phi: M \rightarrow M$ an isometry, and $M := \{p \in M \mid \phi(p) = p\}$. Every connected component of M^ϕ is a closed submanifold of M .

Totally geodesic: A submanifold such that geodesics which start tangent to the submanifold remain there.

Prop 8.22: (M, g) and $\phi: M \rightarrow M$ an isometry. Then the fixed point submanifold M^ϕ is totally geodesic.

8.3 The Riemann Curvature Tensor

Riemann Curvature Tensor: $(0,4)$ -tensor $(\mathfrak{X}(M) \times \mathfrak{X}(M) \times \mathfrak{X}(M) \times \mathfrak{X}(M) \rightarrow C^\infty(M))$
 defined by $\text{Riem}(X, Y, Z, W) := g(R(X, Y)Z, W)$

Prop 8.23: $\forall X, Y, Z, W \in \mathfrak{X}(M)$:

- $\text{Riem}(X, Y, Z, W) = -\text{Riem}(Y, X, Z, W)$
- $\text{Riem}(X, Y, Z, W) + \text{Riem}(Y, Z, X, W) + \text{Riem}(Z, X, Y, W) = 0$
- $\text{Riem}(X, Y, Z, W) = -\text{Riem}(X, Y, W, Z)$
- $\text{Riem}(X, Y, Z, W) = R(Z, W, X, Y)$

Local coords: Riem is determined by its coefficients $R_{ijkl} := \text{Riem}(\partial_i, \partial_j, \partial_k, \partial_l)$, given by
 $R_{ijkl} = g(R(\partial_i, \partial_j)\partial_k, \partial_l) = (R_{ijk}{}^m \partial_m, \partial_l) = R_{ijk}{}^m g_{ml}$.

Have identities: $R_{ijkl} = -R_{jikl} = -R_{ijlk} = R_{klij}$ and $R_{[ijk]l} = 0$.

Ricci tensor : $R_{ij} = R_{ji}$ for Levi-Civita connection

Ricci scalar: $R := g^{ij} R_{ij}$

Dfn 8.27: Einstein tensor of Levi-Civita connection is a symmetric $(0,2)$ -tensor Ein defined by:

$$\text{Ein}(X, Y) := \text{Ric}(X, Y) - \frac{1}{2} g(X, Y) R$$

Local coords: determined by its coefficients: $G_{ij} := R_{ij} - \frac{1}{2} g_{ij} R$

Prop 8.28: $\nabla^j G_{jk} := g^{ij} \nabla_i G_{jk} = 0$ "zero divergence"

Dfn 8.29: (M, g) is said to be Einstein if the Ricci tensor is proportional to the metric:

$$\text{Ric} = \lambda g \quad (\text{or equiv}) \quad R_{ij} = \lambda g_{ij}.$$

for some $\lambda \in \mathbb{R}$. Ricci-flat := $\lambda = 0$.

8.4 Cartan's method of moving frames

Dfn 8.32 (M, g) a Riemannian manifold and $U \subset M$ an open set. A **local orthonormal frame** for M on U is a collection (X_1, \dots, X_n) with $X_i \in \mathcal{X}(U)$ such that $g(X_i, X_j) = \pm \delta_{ij}$ everywhere on U .

$\Rightarrow (X_i)$ are **linearly independent** on U and $\forall p \in U$, $((X_i)_p, \dots, (X_n)_p)$ is an orthonormal frame for the inner product space $(T_p M, g_p)$.

Prop 8.33: (M, g) and $p \in M$. Then \exists local chart (U, x^i) with $p \in U$ and a local orthonormal frame on U .

(M, g) has signature (s, t) : reorder frame so that $\forall p \in U$, (e_1, \dots, e_n) gives an isomorphism $T_p M \cong \mathbb{E}^{s,t}$

notation: $\eta_{ab} := g(e_a, e_b)$

$(\theta^1, \dots, \theta^n) =$ **local coframe canonically dual** to (e_1, \dots, e_n) on U .

$\hookrightarrow \theta^a \in \Omega^1(U)$, $\theta^a(e_b) = \delta_b^a$.

\hookrightarrow then g has local expression $g = \eta_{ab} \theta^a \theta^b$.

Ex 8.34: $e_a^i \theta_i^b = \delta_a^b$, $e_a^i \theta_j^a = \delta_j^i$ for $e_a = e_a^i \partial_i$ and $\theta^a = \theta_i^a dx^i$.

Connection one-forms: $\omega_a^b \in \Omega^1(U)$: $\nabla_x e_a = \omega(x)_a^b e_b$

$\omega_{ab} := \eta_{ac} \omega_c^b$: $\omega_{ab} = -\omega_{ba}$

Prop 8.35: **First structure Equation**: $d\theta^a + \omega_a^b \wedge \theta^b = 0$

Prop 8.36: **Second structure Equation**: $d\omega_a^b + \omega_c^a \wedge \omega_c^b = \Omega_a^b$,

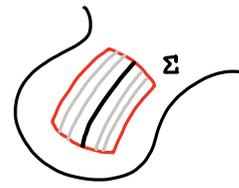
where the **curvature two-form** Ω_a^b is defined by $\Omega(x, y)_a^b e_a = R(x, y)_a^b e_b$

$\hookrightarrow \Omega_{ab} := \Omega_c^b \eta_{ac} \rightarrow \Omega_{ab} = -\Omega_{ba}$

Prop. 8.38. $R_{ijkl} = -\Omega_{ijab} \theta_k^a \theta_l^b$ where $\Omega_{ijab} = \Omega(\partial_i, \partial_j)_{ab}$.

Important identity: $\frac{1}{2} \Omega_{ab} (\theta^a \wedge \theta^b) = -\frac{1}{4} R_{ijkl} (dx^i \wedge dx^j) (dx^k \wedge dx^l)$ } denoting symmetric product of two-forms!

8.5 Geodesic Deviation



o.p.f. of geodesics: a smooth map

$$\gamma: (-\epsilon, \epsilon) \times [0, 1] \rightarrow M, \quad (s, t) \mapsto \gamma_s(t),$$

where for every fixed $s \in (-\epsilon, \epsilon)$, $\gamma_s: [0, 1] \rightarrow M$ is an affinely parametrised geodesic. map parametrizes a 2D surface on M , Σ .

notation: $\dot{\gamma} = \gamma_* \left(\frac{\partial}{\partial t} \right)$, $\gamma' = \gamma_* \left(\frac{\partial}{\partial s} \right)$

Geodesic deviation := acceleration of γ' as we move along the geodesic

$$\frac{D^2}{dt^2} \gamma'$$

Lem 8.42 $\frac{D}{ds} \dot{\gamma} = \frac{D}{dt} \gamma'$

Lem 8.43: Let X be a vector field along Σ . Then $\frac{D}{dt} \frac{D}{ds} X - \frac{D}{ds} \frac{D}{dt} X = R(\dot{\gamma}, \gamma') X$.

Prop 8.44 (Jacobi Equation / Geodesic deviation equation): $\frac{D^2}{dt^2} \gamma' = R(\dot{\gamma}, \gamma') \dot{\gamma}$

\Rightarrow Geodesic deviations determine the curvature tensor.

Lem 8.46: R be the curvature tensor of a torsion free affine connection on a manifold M . Then $\forall X, Y, Z \in \mathfrak{X}(M)$,

$$R(X, Y)Z = \frac{1}{3} \left(R(X, Y)Z + R(Z, Y)X - R(Y, X)Z - R(Z, X)Y \right)$$

Local expression: $R_{ijk}{}^l = \frac{2}{3} (R_{i(jk)}{}^l - R_{j(i)k}{}^l)$

8.6 The Hodge Star Operator

Let (M^4, g) be Lorentzian, $p \in M$. $\exists U \subset M$ s.t. $p \in U$.

Denote the coframe $\theta^a \in \Omega^1(U)$, $g|_U = \eta_{ab} \theta^a \theta^b$, $[\eta_{ab}] = \begin{pmatrix} -1 & 0 \\ 0 & \mathbb{I} \end{pmatrix}$

Dfn: volume form: $\exists \omega \in \Omega^4(M)$ that is nowhere vanishing, $\omega|_U = \theta^0 \wedge \theta^1 \wedge \theta^2 \wedge \theta^3$

Define inner product on $\Omega^1(M)$ by $\langle \theta^a, \theta^b \rangle = \eta^{ab}$, and extend linearly. We can extend this to k -forms by:

$$\langle \theta^a \wedge \theta^b, \theta^c \wedge \theta^d \rangle := \det \begin{pmatrix} \langle \theta^a, \theta^c \rangle & \langle \theta^a, \theta^d \rangle \\ \langle \theta^b, \theta^c \rangle & \langle \theta^b, \theta^d \rangle \end{pmatrix} \Rightarrow \begin{aligned} \langle \theta^0 \wedge \theta^1, \theta^0 \wedge \theta^1 \rangle &= -1 \\ \langle \theta^1 \wedge \theta^2, \theta^1 \wedge \theta^2 \rangle &= 1 \\ &\dots \end{aligned}$$

Dfn: (Hodge star operator): $*$: $\Omega^p(M) \rightarrow \Omega^{4-p}(M)$, $C^\infty(M)$ -linear. e.g. $*$: $\Omega^2(M) \rightarrow \Omega^2(M)$, by

$$\alpha, \beta \in \Omega^2(M), \Rightarrow \alpha \wedge * \beta = \langle \alpha, \beta \rangle \omega$$

Rem: $*(*\alpha) = -\alpha$

Rem: enough to know $*(\theta^a \wedge \theta^b)$

E.g. $(\theta^0 \wedge \theta^1) \wedge *(\theta^0 \wedge \theta^1) = \langle \theta^0 \wedge \theta^1, \theta^0 \wedge \theta^1 \rangle \omega = \langle \theta^0 \wedge \theta^1, \theta^0 \wedge \theta^1 \rangle \omega = -\theta^0 \wedge \theta^1 \wedge \theta^2 \wedge \theta^3$

read off that $*(\theta^0 \wedge \theta^1) = -\theta^2 \wedge \theta^3$

Denote $\theta^{ab} := \theta^a \wedge \theta^b$

α	$*\alpha$
θ^{01}	$-\theta^{23}$
θ^{02}	θ^{13}
θ^{03}	$-\theta^{12}$
θ^{12}	θ^{03}
θ^{13}	$-\theta^{02}$
θ^{23}	θ^{01}

CH 9 Relativity

9.1 Galilean Relativity

Affine space: \mathbb{A}^n (over \mathbb{R}) visualise as affine hyperplane in \mathbb{R}^{n+1} (e.g. last entry equal to 1).

Galilean universe: \mathbb{A}^4 = space of events with:

⊙ **Clock**: $\tau: \mathbb{R}^4 \rightarrow \mathbb{R}$: $\tau(a-b)$ = time between a and b.

↳ **Simultaneous**: $\tau(a-b) = 0$

▮ **Ruler**: $\Delta: \text{Ker } \tau \rightarrow \mathbb{R}$: gives Euclidean distance between simultaneous events.

Affine group: $\text{Aff}(n, \mathbb{R}) \subset \text{GL}(n+1, \mathbb{R})$ which preserves $\mathbb{A}^n \subset \mathbb{R}^{n+1}$.

$$\begin{aligned} &\hookrightarrow \left\{ \begin{pmatrix} A & b \\ 0 & 1 \end{pmatrix} : A \in \text{GL}(n, \mathbb{R}), b \in \mathbb{R}^n \right\} \\ &\hookrightarrow (x, 1), x \mapsto Ax + b \end{aligned}$$

Galilean group: subgroup of $\text{Aff}(n, \mathbb{R})$ preserving clock and ruler.

$$\hookrightarrow \left\{ \begin{pmatrix} R & v & a \\ 0 & 1 & s \end{pmatrix} : v, a \in \mathbb{R}^3, R \in \text{O}(3), s \in \mathbb{R} \right\}$$

breaks down into:

- **Translation**: $\begin{pmatrix} 1 & 0 & a \\ 0 & 1 & s \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ t \\ 1 \end{pmatrix} = \begin{pmatrix} x+a \\ t+s \\ 1 \end{pmatrix}$
- **Galilean boost**: $\begin{pmatrix} 1 & v & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ t \\ 1 \end{pmatrix} = \begin{pmatrix} x+tv \\ t \\ 1 \end{pmatrix}$ (v = "velocity")
- **Reorientation of axes**: $\begin{pmatrix} R & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ t \\ 1 \end{pmatrix} = \begin{pmatrix} Rx \\ t \\ 1 \end{pmatrix}$

Principal idea of Galilean Relativity: "inertial coordinates": free particles move in straight lines ("world line of particle")

9.2. Special Relativity

Laplacian: $\Delta := \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$

d'Alembertian: $\square := \frac{\partial}{\partial t^2} - \frac{\partial^2}{\partial x^2} - \frac{\partial^2}{\partial y^2} - \frac{\partial^2}{\partial z^2}$

Relative to Minkowski coords, $\square = -\eta_{\mu\nu} \partial_\mu \partial_\nu$

Minkowski Universe: $\mathbb{A}^4 =$ space of spacetime events + "proper distance" between two spacetime events $a, b \in \mathbb{A}^4$:

$$\Delta(a, b) = \eta(b - a, b - a)$$

$\eta =$ Minkowski inner product defined on \mathbb{R}^4

$b - a = (t, x, y, z) \implies \Delta(a, b) = -t^2 + x^2 + y^2 + z^2.$

-ve prop. dist =: proper time.

Poincaré group: subgroup of $\text{Aff}(4, \mathbb{R})$ preserving proper distance.

$$\hookrightarrow \left\{ \begin{pmatrix} A & v \\ 0 & 1 \end{pmatrix} : A \in O(3,1), v \in \mathbb{R}^4 \right\}$$

$$\eta = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}.$$

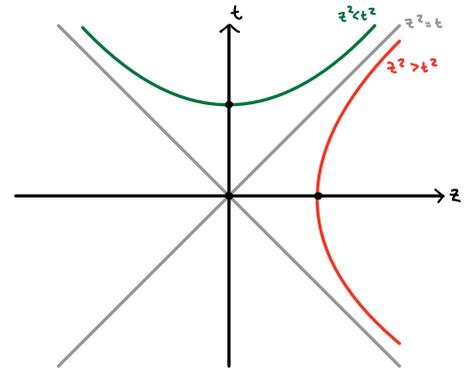
Where Lorentz group: $O(3,1) = \{ A \in GL(4, \mathbb{R}) : A^T \eta A = \eta \}$,

Lorentz boost: $(t, x, y, z) \mapsto (t', x, y, z')$

↓ Where by linearity $z' = \alpha z + \beta t$, $t' = \gamma z + \delta t$.

Reparametrise: $z' = z \cosh \zeta + ct \sinh \zeta$,

$t' = \frac{1}{c} z \sinh \zeta + t \cosh \zeta.$



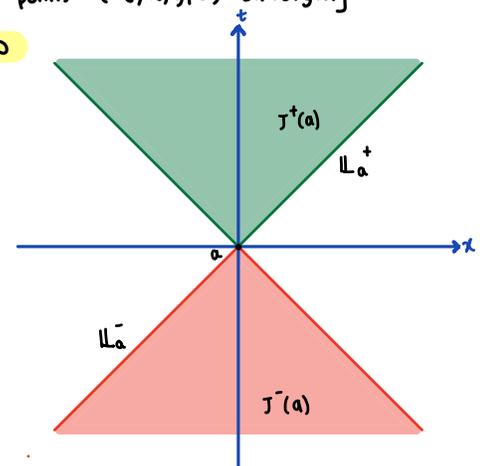
Nonzero Poincaré-invariant clock $\tau: \mathbb{R}^4 \rightarrow \mathbb{R}$ in Minkowski spacetime.

$\mathbb{L}_a =$ lightcone based at a : if $a = (t_0, x_0, y_0, z_0)$, then \mathbb{L}_a consists of points (t, x, y, z) satisfying

$$-(t - t_0)^2 + (x - x_0)^2 + (y - y_0)^2 + (z - z_0)^2 = 0$$

\mathbb{L}_a^+ = future lightcone $t > t_0$

\mathbb{L}_a^- = past lightcone $t < t_0$.



a, b are causally related if $\Delta(a, b) \leq 0$.

a, b are spacelike separated if $\Delta(a, b) > 0$

$J^+(a) =$ Causal future: points at a non+ proper dist. to a

$J^-(a) =$ Causal past: points at a non- proper dist. to a

Relativistic equations: PDE's for functions on Minkowski spacetime that are invariant under \square

Maxwell's equations in vacuo: $dF = 0$ and $d \star F = 0$

9.3 General Relativity

Postulates of GR:

- 1 Spacetime is a connected, 4D Lorentzian manifold (M, g)
- 2 Free particles in the spacetime follow timelike or null geodesics of the Levi-Civita connection

Minimal coupling: go from SR to GR

- 3 The distribution of matter (including radiation) is described by the energy-momentum tensor T , a symmetric $(0,2)$ -tensor which has zero divergence $\nabla^\mu T_{\mu\nu} = 0$.

Exm 9.18: Maxwell's equations: let $F \in \Omega^2(M)$, $F = \frac{1}{2} F_{\mu\nu} dx^\mu \wedge dx^\nu$.

Equations are: $dF = 0$ $g^{\mu\nu} \nabla_\mu F_{\nu\rho} = 0$

Energy-momentum tensor of F : $T_{\mu\nu} = g^{\rho\sigma} F_{\mu\rho} F_{\nu\sigma} - \frac{1}{4} g_{\mu\nu} g^{\rho\tau} g^{\sigma\chi} F_{\rho\sigma} F_{\tau\chi}$

- 4 The curvature of the spacetime is related to the energy-momentum tensor via the Einstein Equation

$$E_{in} = 8\pi G T$$

or relative to a local chart,

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Taking trace: Einstein equation then becomes $Ric = \Lambda g$, where Λ is called the "cosmological constant".

CH 10 Exact Solutions

10.1 Isometric actions of Lie Groups

10.1.1 Group actions on Riemannian manifolds $G :=$ Lie group, M a manifold.

Left action of G on M is a smooth map $\alpha: G \times M \rightarrow M$, $(g, p) \mapsto g \cdot p$ satisfying:
 1 $\forall p \in M, g_1, g_2 \in G, g_1 \cdot (g_2 \cdot p) = (g_1 g_2) \cdot p$ 2 $\forall p \in M, e \cdot p = p$

$\forall g \in G$, get an associated diffeo $\phi_g: M \rightarrow M, p \mapsto g \cdot p$. Note: $\phi_g \circ \phi_h = \phi_{gh}$

- A left action is
- **Effective:** $\ker \phi = \{e\}$
 - **Transitive:** $\forall p, q \in M, \exists g$ s.t. $p = g \cdot q$, equivalently $G \cdot p = M, \forall p$

Lem 10.1: G compact. Then $G \cdot p \subset M$ is an embedded subman. of $M, \forall p$.

Fundamental vector field: $\mathfrak{g} = T_e G$, and $X \in \mathfrak{g}$. Then X defines a vector field \tilde{X} on M as follows:
 Let $\gamma: (-\epsilon, \epsilon) \rightarrow G$ be a curve with $\gamma(0) = e$ and $\gamma'(0) = X \in \mathfrak{g}$. Given $p \in M$, let $c: (-\epsilon, \epsilon) \rightarrow M$ be the curve in M s.t. $c(t) = \gamma(-t) \cdot p$. Then $\tilde{X}_p := c'(0) \in T_p M: \forall f \in C^\infty(M)$,

$$(\tilde{X}f)(p) = \left. \frac{d}{dt} f(\gamma(-t) \cdot p) \right|_{t=0}$$

Lem 10.3: $\forall X, Y \in \mathfrak{g}, [\tilde{X}, \tilde{Y}] = \widetilde{[X, Y]}$ $G_p :=$ Stabilizer of p

Thm 10.5 (Orbit-Stabilizer Theorem) Let G act smoothly on M and $p \in M$. \exists a G -equivariant diffeo

$$\varphi: G \cdot p \rightarrow G/G_p$$

defined by $\varphi(p) = G_p$ and $\varphi(g \cdot p) = gG_p$.

Clearly G -equivariant: $\varphi(g \cdot p) = g \cdot \varphi(p)$.

10.1.2: Homogeneous Riemannian Manifolds

Dfn: (M, g) is **homogeneous** if it admits a transitive isometric action of a Lie Group.

I.e. \exists Lie Group G s.t. $\forall p \in M, G \cdot p = M$ and G preserves the metric.

Ex: Minkowski is hom: translations act transitively.

Thm 10.8 (Frobenius Reciprocity): 1-1 Correspondence between G_p -invariant tensors on $T_p M$ and G -invariant tensor fields on M .

10.2 Spherical Symmetry

$G :=$ Lie group $SO(3)$, Orbits = $r \in [0, \infty)$

Dfn 10.9: (Spherical symmetry) (M, g) a 4D Lorentzian man. on which G acts by isometries.

(M, g) is **spherically symmetric** if the generic G -orbits are 2-spheres.

Prop 10.14: (M, g) a spherically symmetric spacetime. Then either it is isometric to the Riemannian product $(k, g_k) \times (s^2, r_0 g_s^2)$ of a Lorentzian two-dimensional manifold and a round sphere of radius r_0 , or else \exists chart with local coordinates (τ, r, θ, ϕ) , relative to which the metric takes the form

$$g = g_{\tau\tau}(\tau, r) d\tau^2 + g_{rr}(\tau, r) dr^2 + r^2(d\theta^2 + \sin^2\theta d\phi^2).$$

10.3 Schwarzschild Metric

Deriving a Ricci-flat metric.

10.3.1 Deriving the Schwarzschild Metric

Dfn 10.15: (M, g) is called...

Stationary: if it has a timelike $(g(X, X) < 0)$ Killing vector, X

Static: stationary and the one form $\theta = X^\flat$ canonically dual to X satisfies $\theta \wedge d\theta = 0$.

↳ X is then "hypersurface orthogonal"

E.g. Minkowski, $X = \frac{\partial}{\partial t}$ (clearly $g(X, X) = -1 < 0$).

Ex 10.17: $\theta \in \Omega^1(M)$, $\theta \wedge d\theta = 0$

⇒ $\exists \omega \in \Omega^1(M)$ s.t. $d\theta = \theta \wedge \omega$,

⇒ $\gamma, M \in \mathcal{X}(M)$ s.t. $\theta(\gamma) = \theta(z) = 0$, ⇒ $\theta([\gamma, z]) = 0$

Schwarzschild metric: $g = -\left(1 - \frac{2GM}{r}\right) dt^2 + \left(1 - \frac{2GM}{r}\right)^{-1} dr^2 + r^2 d\theta^2 + r^2 \sin^2\theta d\phi^2$

$r \rightarrow \infty$, $g \rightarrow$ Minkowski

↳ " g is asymptotically flat"

10.3.2: Singularities of the Schwarzschild metric.

Singularities:

$r = r_s := 2GM$
coordinate Singularity

$r = 0$

physical Singularity

Kretschmann Scalar: $K := \frac{1}{4} R^{\mu\nu\rho\sigma} R_{\mu\nu\rho\sigma}$

Prop 10.20: The Kretschmann Scalar for the Schwarzschild metric is given by

$$K = \frac{12M^2}{r^6}$$

⇒ $r=0$ is actually a physical Singularity.

10.4: The Schwarzschild Black Hole

10.4.1: The Rindler Wedge:

2D Lorentzian metric: $g = dx^2 - x^2 dt^2$, $(z, t) \in (0, \infty) \times \mathbb{R}$

10.4.2: The Kruskal Extension

Concentrate on (r, t) part of Schwarzschild metric: $g = -\left(1 - \frac{2GM}{r}\right) dt^2 + \left(1 - \frac{2GM}{r}\right)^{-1} dr^2$

rewriting using change of coordinates: total metric: $g = 16M^2 \frac{1}{r} e^{-r/2M} (dx^2 - dt^2) + r(d\theta^2 + \sin^2\theta d\phi^2)$

Dfn 10.24: (M, g) be a spacetime and $j: \Sigma \rightarrow M$ an

embedded hypersurface. We can pull back g via j

onto Σ to give a symmetric $(0, 2)$ -tensor field

$j^*g \in \Gamma(\mathcal{O}^2 T^* \Sigma)$. We say that Σ is...

- **spacelike:** j^*g is positive-definite
- **timelike:** j^*g is a Lorentzian metric
- **null:** j^*g is degenerate

